Measuring the Effect of LQC in the CMB Data (tachyon warm model)

Vahid Kamali

BASU and IPM-Iran

Paris, 31 Aug, 2017

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1 Introduction

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Introduction

- 2 Why Warm Inflation?
 - Energy density evolution
 - Evolution of warm inflation
 - Condition of Warm inflation

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Introduction

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 - Pressure and energy density

Introduction

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4 LQC

Introduction

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 - Evolution of warm inflation
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- 3 Why Tachyon Field?
 - Pressure and energy density
- 4 LQC
- 5 Results

Cold Inflation

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Inflation

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Inflation

Standard Model of inflation

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Inflation

- Standard Model of inflation
- Slow-roll approximation

Cold Inflation

Inflation

- Standard Model of inflation
- Slow-roll approximation
- Reheating

Standard inflation

condition

We know from G.R in order to realize inflation, it requires an equation of state $P < -\rho/3$, thus a substance with negative pressure that scalar fields can provide such an equation of state.

$$\rho(\phi) = \frac{1}{2}\dot{\phi}^2 + V(\phi) \qquad P(\phi) = \frac{1}{2}\dot{\phi}^2 - V(\phi)$$
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$$\rho(\phi) = \frac{1}{2}\dot{\phi}^2 + V(\phi) \qquad P(\phi) = \frac{1}{2}\dot{\phi}^2 - V(\phi)$$
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Evolution of the scalar field

Related equations

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Evolution of the scalar field

Related equations

Energy density conservation

$$\dot{\rho} + 3H(\rho + P) = 0 \; , \qquad \qquad$$

(2)

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Evolution of the scalar field

Related equations

Energy density conservation

$$\dot{\rho} + 3H(\rho + P) = 0 , \qquad (2)$$

The equations describing the evolution scalar field

$$\ddot{\phi} + 3H\dot{\phi} + rac{dV}{d\phi} = 0 \;,\; H^2 = rac{8\pi G}{3}
ho$$

where H is Hubble parameter

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The other dynamical realization of inflation is warm inflation. In this picture, similar to cold inflation, the scalar inflaton field must be potential energy dominated to realize inflation

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Conservation equation

$$\begin{split} \dot{\rho}_{\phi} + 3H(\rho_{\phi} + P_{\phi}) &= -\Gamma \dot{\phi}^2 \\ \dot{\rho}_{\gamma} + 4H\rho_{\gamma} &= \Gamma \dot{\phi}^2; \end{split}$$

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Conservation equation

$$\dot{
ho}_{\phi} + 3H(
ho_{\phi} + P_{\phi}) = -\Gamma\dot{\phi}^2$$

 $\dot{
ho}_{\gamma} + 4H
ho_{\gamma} = \Gamma\dot{\phi}^2;$

Equation of motion of warm inflation

$$\ddot{\phi} + (3H + \Gamma)\dot{\phi} + \frac{dV}{d\phi} = 0$$

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Perturbation of the scalar field

$$\delta \phi_{warm}^2 \sim \sqrt{HT} \qquad \qquad \delta \phi_{cold} \sim H;$$

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$$\delta \phi^2_{warm} \sim \sqrt{HT} \qquad \qquad \delta \phi_{cold} \sim H$$

In warm inflation model we have T > H.

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Perturbation of warm inflation, first approximation

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Perturbation of warm inflation, first approximation

Power spectrum

$$\mathcal{P}_R \sim \frac{H^2}{\dot{\phi}^2} \delta \phi^2;$$

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Perturbation of warm inflation, first approximation

Power spectrum

$$\mathcal{P}_R \sim \frac{H^2}{\dot{\phi}^2} \delta \phi^2;$$

Tensor-scalar ratio

$$r = \left(\frac{H}{T}\right) 16\epsilon;$$

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Perturbation of warm inflation, first approximation

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Tensor-scalar ratio

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Spectral index

$$n_s = 1 + \frac{d\ln \mathcal{P}_R}{d\ln k};$$

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Perturbation of warm inflation, first approximation

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Tensor-scalar ratio

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$$n_s = 1 + \frac{d\ln \mathcal{P}_R}{d\ln k};$$

We can present these parameters for our model

Image: Image:

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From the dynamical viewpoint one may obtain the equations of motion of tachyon fields using a special Lagrangian

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- From the dynamical viewpoint one may obtain the equations of motion of tachyon fields using a special Lagrangian
 - Lagrangian

From the dynamical viewpoint one may obtain the equations of motion of tachyon fields using a special Lagrangian

1 Lagrangian

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$$L = \sqrt{-g} \left[\frac{R}{16\pi G} - V(\phi) \sqrt{1 - g^{\mu\nu} \partial_{\mu} \phi \partial_{\nu} \phi} \right] .$$
(9)

2 Stress-energy tensor

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From the dynamical viewpoint one may obtain the equations of motion of tachyon fields using a special Lagrangian

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2 Stress-energy tensor

$$T^{\mu}_{\nu} = \frac{\partial L}{\partial(\partial_{\mu}\phi)} \partial_{\nu}\phi - g^{\mu}_{\nu}L = \text{diag}(-\rho_{\phi}, p_{\phi}, p_{\phi}, p_{\phi})$$
(10)

equation

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equation

Pressure and energy density of tachyonic fluid in FRW universe we have

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equation

Pressure and energy density of tachyonic fluid in FRW universe we have

$$\rho_{\phi} = \frac{V(\phi)}{\sqrt{1 - \dot{\phi}^2}} \qquad P_{\phi} = -V(\phi)\sqrt{1 - \dot{\phi}^2}$$
(11)

slow-roll limit

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slow-roll limit

$$\dot{\phi}^2 \ll 1 \quad \Rightarrow P = -\rho$$

Notice, that tachyon potentials have the following two properties: the maximum of the potential occurs when $\phi \to 0$ while the corresponding minimum takes place when $\phi \to \infty$.



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LQC

LQC is LQG limit in FRW universe which modified Friedmann equation in a sense that the singularity is removed.

Modification of first Friedmann equation

$$H^{2} = \frac{8\pi G}{3}\rho(1 - \frac{\rho}{\rho_{c}})$$
 (12)

where ρ_c is constant

Warm LQC

$$\rho = \rho_{\phi} + \rho_{\gamma} \qquad P = P_{\phi} + P_{\gamma} \tag{13}$$

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Results

Summary of results

we estimate analytically the slow-roll parameters and we compare our predictions with those of other inflationary models as well as we test the performance of warm LQC inflation against the observational data.



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Thanks for your attention

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