

A new formalism for cosmological observables (without coordinates)

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Observation and tetrads

- *What do we measure?*

Spectrum, flux, intensity, polarization, angular position, shape and size in the sky of light beams

- *Where does this information lie?*

The momentum vector and polarization tensor of photons

⇒ We measure the components of tensors at $\mathcal{O} \in \mathcal{M}$

⇒ w.r.t. a basis of $T_{\mathcal{O}}\mathcal{M}$ privileged by the apparatus

- This basis e_a^μ is orthonormal, i.e. a tetrad

$$g_{\mu\nu}(\mathcal{O}) e_a^\mu e_b^\nu = \eta_{ab} \quad \text{and} \quad e_0^\mu \equiv u_{\mathcal{O}}^\mu$$

and we measure $k^a(\mathcal{O}) := e_{\mu}^a k^\mu(\mathcal{O}) \equiv \omega_{\mathcal{O}} (1, -n_{\mathcal{O}}^i)$

- In the metric formalism the only available frames are ∂_μ

$$g(\mathcal{O})(\partial_\mu, \partial_\nu) \equiv g_{\mu\nu}(\mathcal{O}) \neq \eta_{\mu\nu}$$

especially in the usual gauges (Newtonian and synchronous)

\Rightarrow the components k^μ are *not* the measured ones

- We do have, however, $u_{\mathcal{O}}^\mu$ and $u_{\mathcal{S}}^\mu$

\Rightarrow we can define *some* observables such as redshift

$$1 + z := \frac{(u_\mu k^\mu)_{\mathcal{S}}}{(u_\mu k^\mu)_{\mathcal{O}}}$$

- But we still lack a spatial orthonormal frame $e_{i=1,2,3}^\mu$

\Rightarrow we cannot define angle-dependent observables

- Solution 1: choose coordinates such that $g_{\mu\nu}(\mathcal{O}) = \eta_{\mu\nu}$
 - \Rightarrow fixes the gauge freedom in the vicinity of \mathcal{O}
 - \Rightarrow no matching with standard gauges on the rest of \mathcal{M}

- Solution 2: invoke a tetrad e_a^μ at \mathcal{O} only
and build it out of $g_{\mu\nu}(\mathcal{O})$ and $u_{\mathcal{O}}^\mu$
 - \Rightarrow 16 components in e_a^μ versus 13 in $g_{\mu\nu}(\mathcal{O})$ and $u_{\mathcal{O}}^\mu$
 - \Rightarrow partial gauge-fixing – usual choices break diffs near \mathcal{O}
 - \Rightarrow again, no freedom to go to standard gauges in \mathcal{M}

Our solution: tetrad formulation of GR

- Replace the metric information by a tetrad *field* $e_a^\mu(x)$

$$g_{\mu\nu}(x) = \eta_{ab} e_\mu^a(x) e_\nu^b(x)$$

- Advantage 1:

Unifies observer and sources into an “observer family”, e.g.

$$u_{\mathcal{O}}^\mu \equiv e_0^\mu(\mathcal{O}), \quad u_{\mathcal{S}}^\mu \equiv e_0^\mu(\mathcal{S})$$

$\Rightarrow e_i^\mu$ spans the rest-frame directions by construction

- Advantage 2:

Work directly with $T_{b_1 \dots b_m}^{a_1 \dots a_n}(x)$ on \mathcal{M} (diff scalars)

$\Rightarrow T_{b_1 \dots b_m}^{a_1 \dots a_n}(\mathcal{O})$ are the observed quantities

\Rightarrow explicitly coordinate-independent formalism

- Advantage 3:

New gauge symmetry: local Lorentz transformations (LLT)

$$\tilde{e}_\mu^a(x) = \Lambda^a_b(x) e_\mu^b(x) \quad \Rightarrow \quad \tilde{g}_{\mu\nu}(x) = g_{\mu\nu}(x)$$

▶ 16 tetrad components \rightarrow 10 metric components

▶ LLT = “observer family transformations”

k^a measured by $e_a^\mu \rightarrow \tilde{k}^a$ measured by \tilde{e}_a^μ

\Rightarrow local $SO(1,3)$ allows to reach all possible observer families

▶ Conceptual transparency

Symmetry \Rightarrow observer-*independent* physics

Relativity \Rightarrow observer-*dependent* measurements

\Rightarrow Coordinates are no longer related to these concepts

- Advantage 4:
Covariance under both gauge transformations
⇒ Freedom to choose diff gauges (e.g. Newtonian)
⇒ LLT gauge amounts to observer dynamics (e.g. free-fall)
- Advantage 5:
Unifies gravitational and observer information:
Gravity = internal norm of the tetrad (metric, LLT-scalar)
Observer = orientation of the tetrad (LLT-vector)
- Advantage 6:
Symmetry-based definition of observables $C(\mathcal{S}, \mathcal{O})$
 - ▶ Invariant under coordinate transformations
 - ▶ Can only depend on $\Lambda(\mathcal{O})$ and $\Lambda(\mathcal{S})$ under LLTs

Example: the angular diameter distance

- Deviation of the light-like geodesic in the tetrad basis

$$d^a := e_{\mu}^a(\gamma) \delta_{\text{dev}} \gamma^{\mu} \quad d_{\perp}^i := d^i - n^i n^j d^j$$

- Deviation of the observed position in the sky

$$m_{\mathcal{O}}^i := \delta_{\text{dev}} n_{\mathcal{O}}^i \sim d_{\mathcal{O}}^i \quad m_{\mathcal{O}}^i n_{\mathcal{O}}^i \equiv 0$$

- Physical area of the source (visible face) in its rest-frame

$$dA_S := \varepsilon^{ijk} n_S^i d_S^j \tilde{d}_S^k$$

- Observed solid angle

$$d\Omega_{\mathcal{O}} := \varepsilon^{ijk} n_{\mathcal{O}}^i m_{\mathcal{O}}^j \tilde{m}_{\mathcal{O}}^k$$

Observation implies

$$d_{\mathcal{O}}^i = 0$$

\Rightarrow geodesic deviation solution takes the form

$$d_{\perp, \mathcal{S}}^i = \mathcal{J}_{\mathcal{S}, \mathcal{O}}^{ij} m_{\mathcal{O}}^j$$

where \mathcal{J} is the lensing Jacobi map obeying

$$n_{\mathcal{S}}^i \mathcal{J}_{\mathcal{S}, \mathcal{O}}^{ij} \equiv 0 \quad \mathcal{J}_{\mathcal{S}, \mathcal{O}}^{ij} n_{\mathcal{O}}^j \equiv 0$$

We then have

$$dA_{\mathcal{S}} \equiv D_{\mathcal{S}, \mathcal{O}}^2 d\Omega_{\mathcal{O}}$$

$$D_{\mathcal{S}, \mathcal{O}}^2 = \det_{\perp} \mathcal{J} \equiv \frac{1}{2} (n_{\mathcal{S}}^i \varepsilon^{ijk}) (n_{\mathcal{O}}^k \varepsilon^{klm}) \mathcal{J}^{jl} \mathcal{J}^{km}$$

No coordinates involved.

Thank you!