

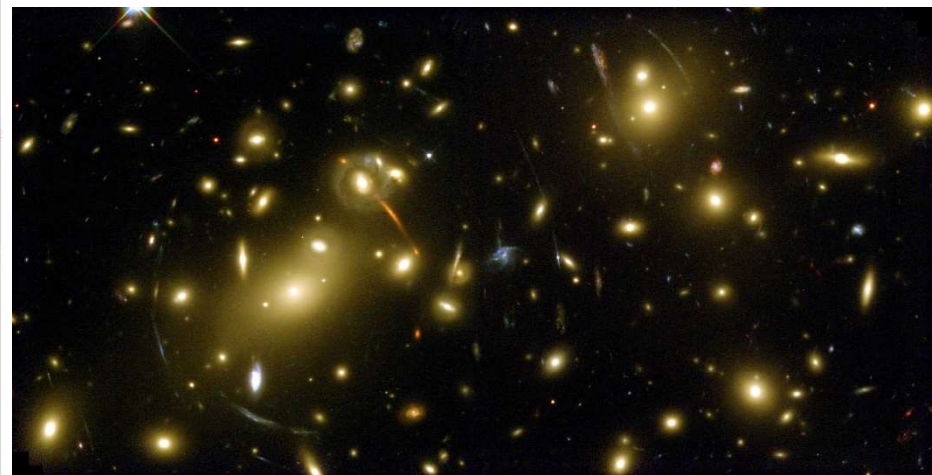
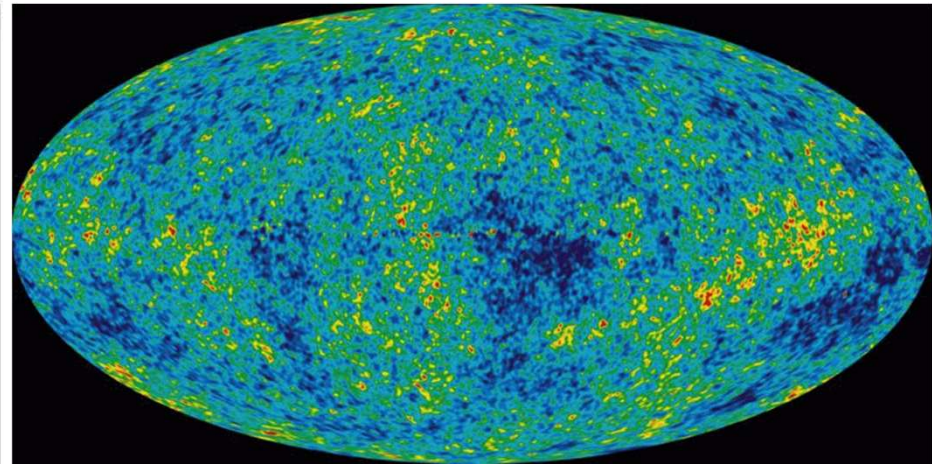
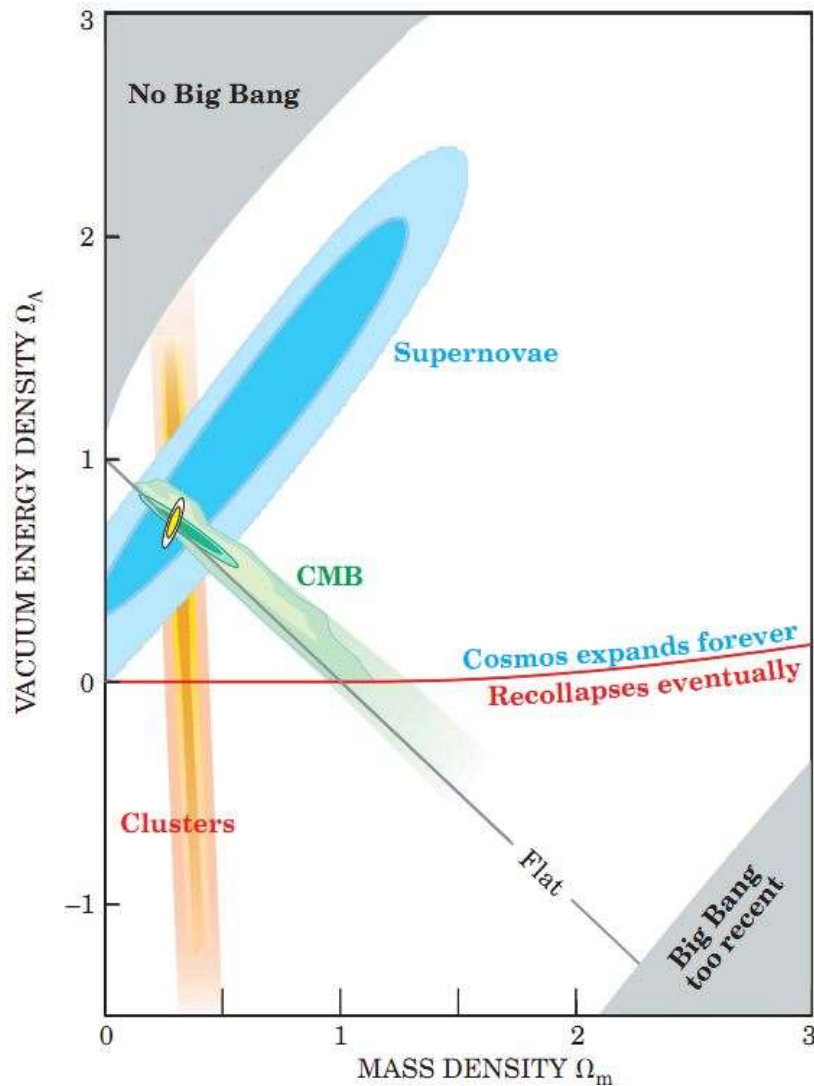
# Gauss-Bonnet Coupled Quintessential Inflation

arXiv: 1707.06839

28th August 2017  
Charlotte Owen

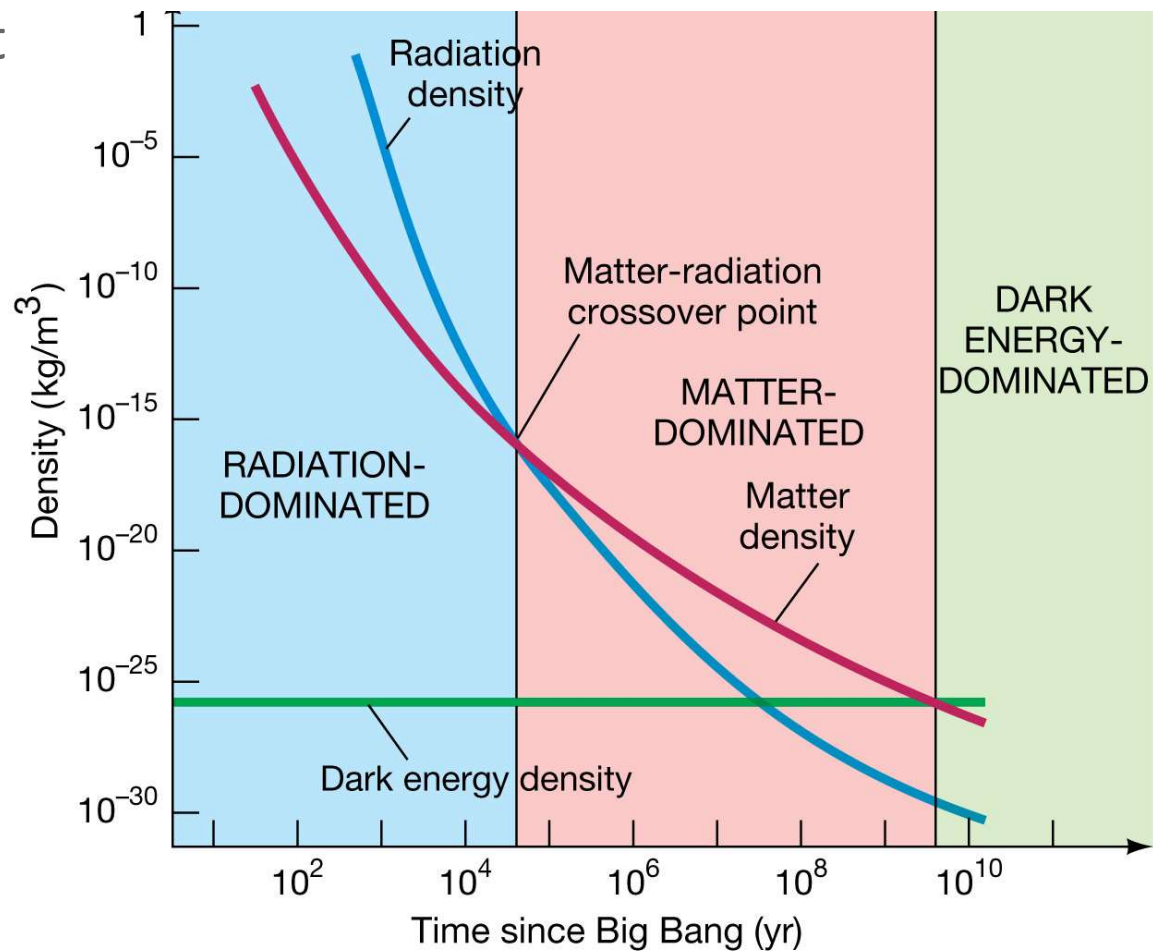
Work done under the supervision of Dr. Konstantinos Dimopoulos, in collaboration with Prof. Carsten van de Bruck and Chris Longden from Sheffield University.

# Accelerating Expansion



# Cosmological Constant Model

- Cosmological Constant
  - energy density of empty vacuum with negative pressure
- $\rho_{\Lambda} \approx (10^{-3} \text{ eV})^4$
- Equation of state parameter:  $\omega = -1$
- Problems with CC:
  - Cosmological Constant Problem
  - Fine-tuning
  - Future Horizons




# Quintessence


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- Dynamical scalar field, varying in space and/or time
- Dominates the energy density of the Universe at the present time

$$\omega \leq -\frac{1}{3}$$

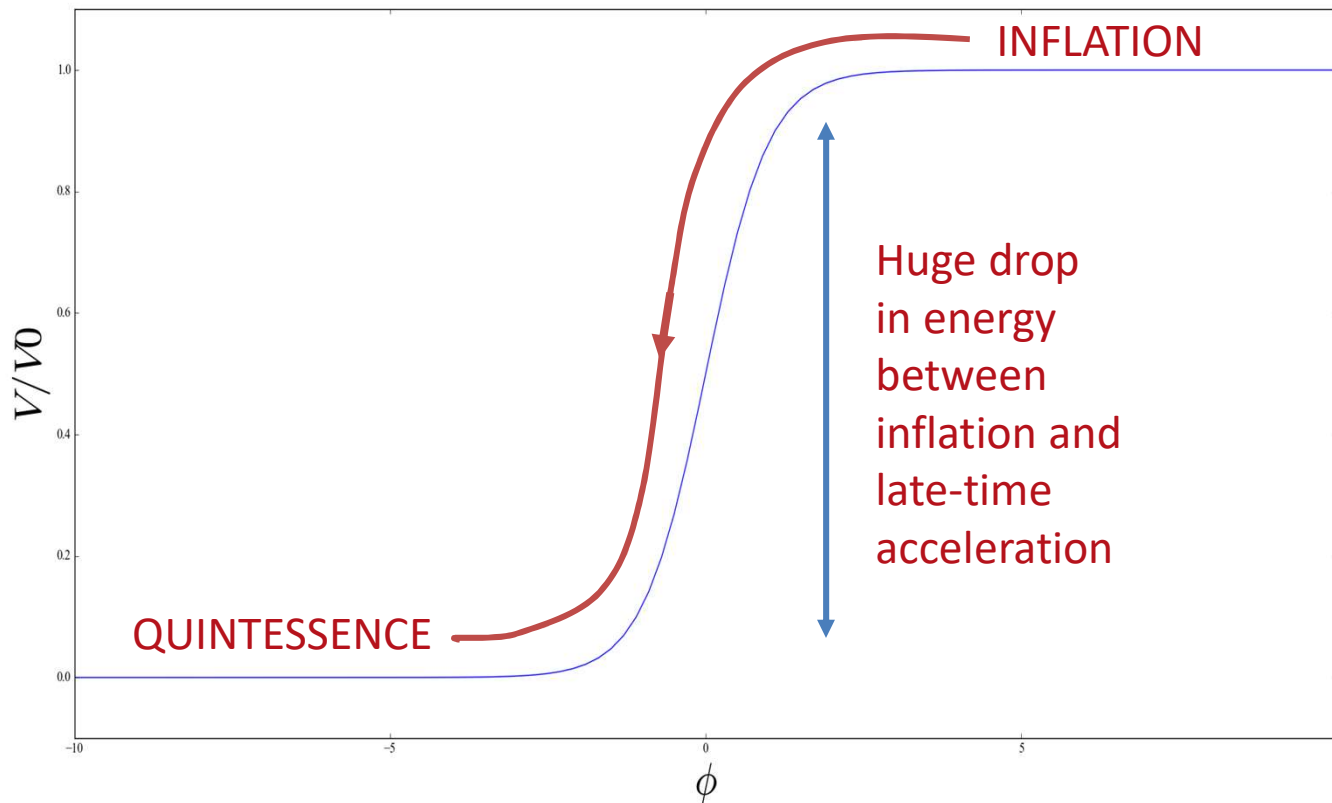
- Ratra & Peebles, 1988 (before discovery of Dark Energy!)
  - Problems with Quintessence:
    - Fine tuning problem persists (initial conditions)
    - Fifth force problem
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# Quintessential *Inflation*

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- **Links** the two observed phases of accelerated expansion – inflation and dark energy, with the **same mechanism**.
  - Considerations:
    - Persisting fine-tuning problem but initial conditions fixed.
    - Persisting fifth force problem.
    - Very large difference between plateaus – must work at two very different energy scales whilst maintaining flat plateaus.
    - Avoid super-planckian field values.
    - Should not feature  $\Lambda$  scale:  $\rho_\Lambda \approx (10^{-3} \text{ eV})^4$
- 

# Tanh as a Prototype Potential

$$V(\phi) = \frac{V_0}{2} \left( 1 + \tanh \left[ p \frac{(\phi - \phi_c)}{m_P} \right] \right)$$



- Two plateaus
- Rolls super-Planckian when the exponential is very steep
- Radiative corrections threaten flatness of late-time plateau
- How can we fix this?

# Introducing a Gauss-Bonnet Term

$$S = \int d^4x \sqrt{-g} \left[ \frac{m_P^2}{2} R - \frac{1}{2} (\partial\phi)^2 - \frac{m_P^2}{2} G(\phi) E - V(\phi) \right]$$

**Gauss-Bonnet:**  $E = R^2 - 4R^{\mu\nu}R_{\mu\nu} + R^{\mu\nu\rho\sigma}R_{\mu\nu\rho\sigma}$

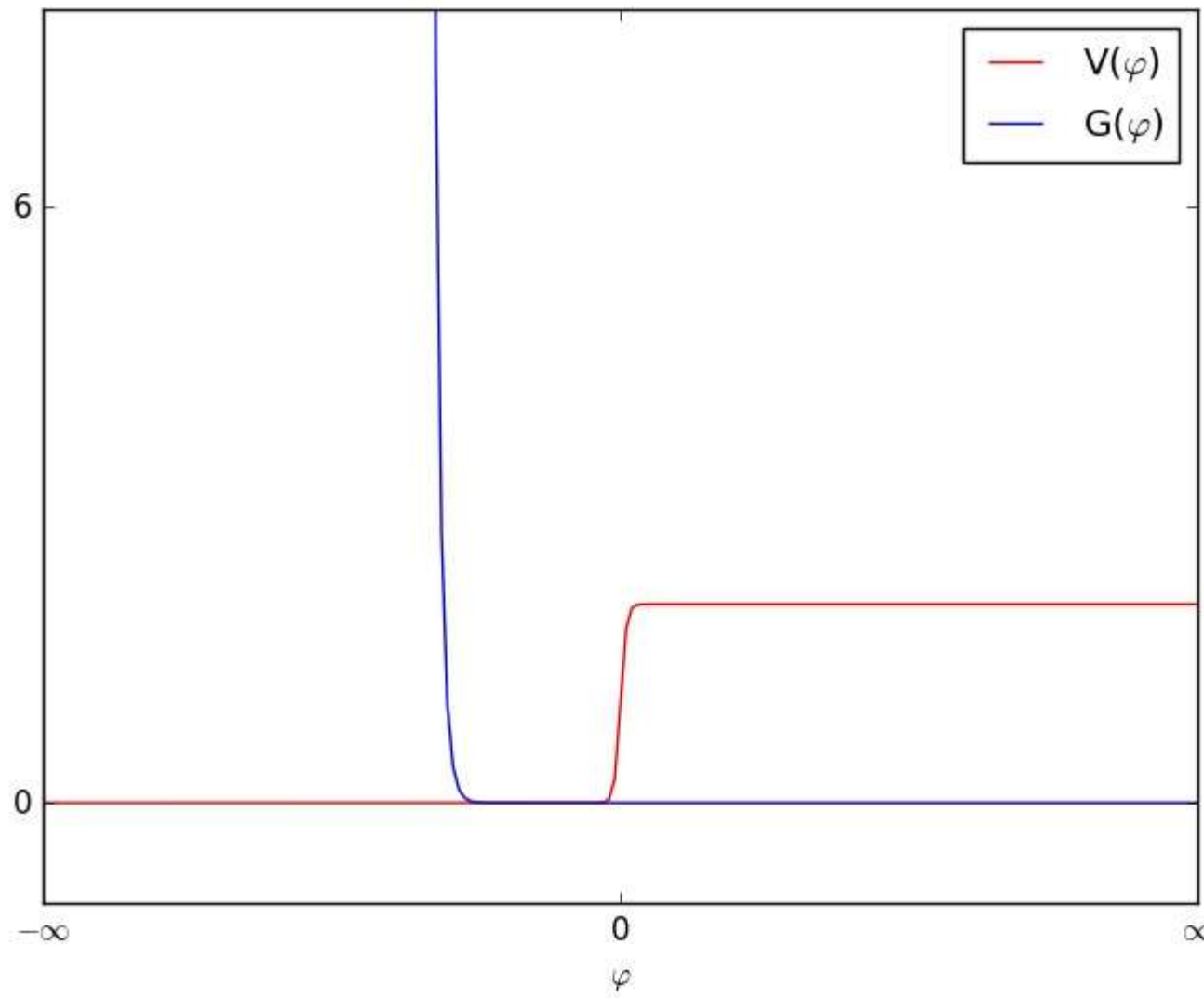
- Total derivative
- No ghosts
- No effect on equations of motion

**Coupling to the Gauss-Bonnet:**  $G(\phi) = G_0 e^{\frac{-q\phi}{m_P}}$

- Varies with  $\phi$
- Affects equations of motion
- Suppressed until late times
- Freezes the inflaton's motion



# GB Introduces Minimum to Effective Potential



- Minimum in effective potential
- Arises from non-constant coupling to GB term
- We can control its location



# Inflationary Observables

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$$n_s \approx 1 - \frac{2}{N}$$


$$r \approx \frac{2}{p^2 N^2}$$

For  $N = 50$  (60):

$$n_s \approx 0.960 \text{ (0.967)}$$

$$r_{\max} \approx 0.0008 \text{ (0.0006)}$$

Planck  $1\sigma$  ( $2\sigma$ ):  $n_s = 0.968 \pm 0.006$  (0.01)



# Model Constraints

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Remain sub-Planckian:

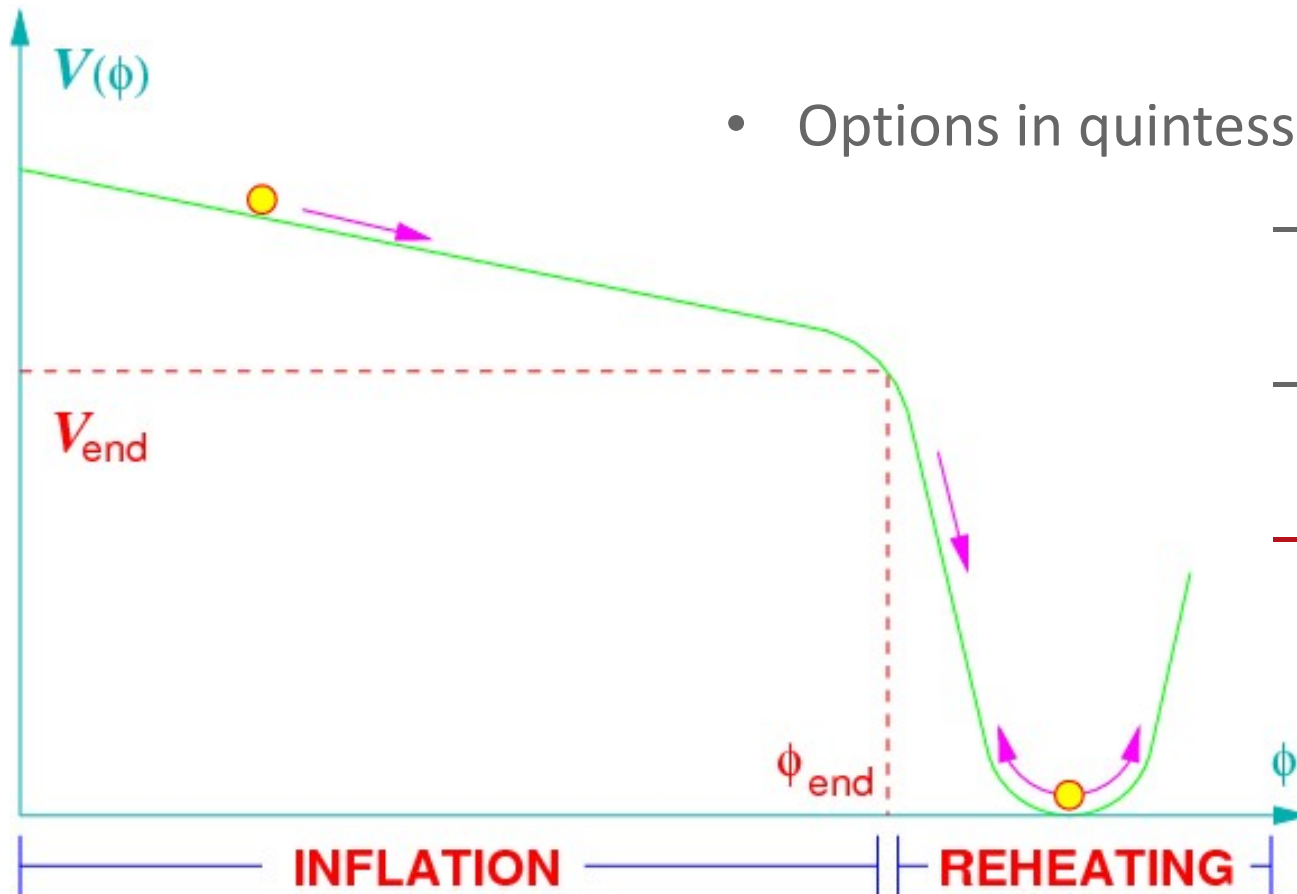
$$\phi < m_P$$

Satisfy dark energy observations:

$$\rho(\phi_0) = \rho_{DE} = 10^{-120} m_P^4$$

# Reheating in Quintessential Inflation

Models of quintessential inflation have non-oscillatory behaviour. Cannot reheat the Universe in the traditional perturbative way.



- Options in quintessential inflation:

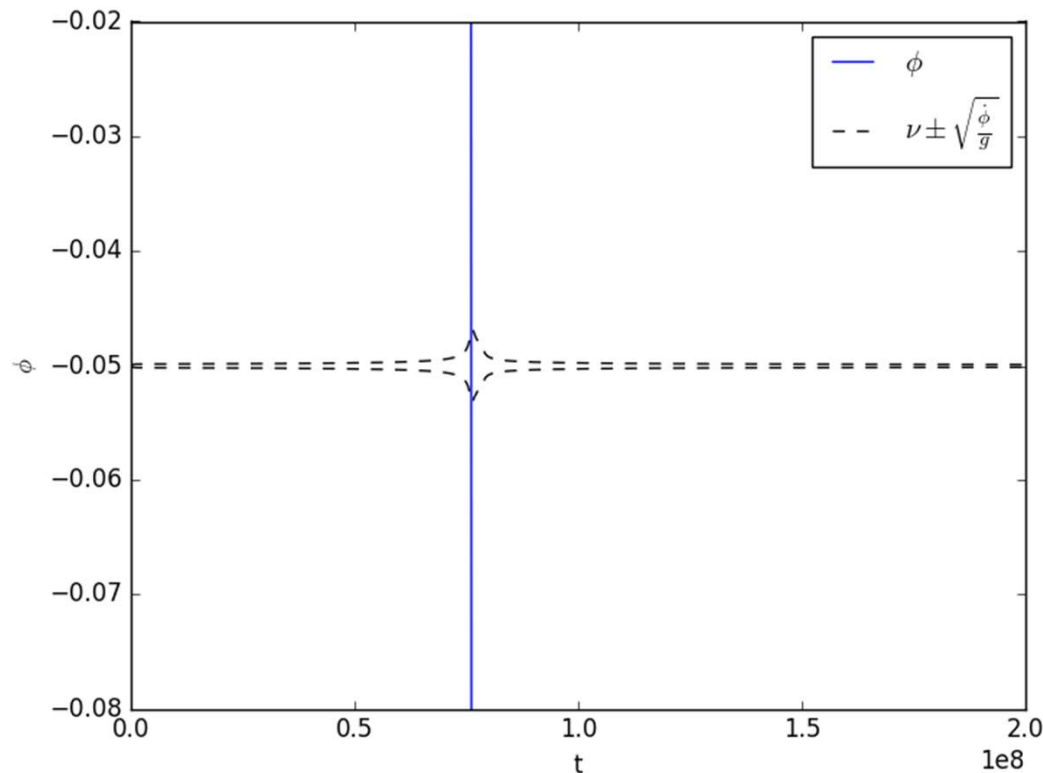
- Gravitational reheating
- Curvaton reheating
- Instant preheating

# Instant Preheating

Felder, Kofman & Linde, Phys. Rev. D 59 (1999)

$$\mathcal{L} = -\frac{1}{2} g^2 \chi^2 |\phi - v|^2 - h \bar{\varphi} \varphi \chi$$

Adiabaticity constraint violated, particle production:  $|\dot{m}_\chi| \gg m_\chi^2$



- Region of particle production:

$$\phi = v \pm \sqrt{\frac{|\dot{\phi}|}{g}}$$

- Shift symmetry allows us to incorporate instant preheating in to a range of model parameters

# Correct History of Universe Evolution

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For radiation domination:  $\rho(\chi^*) > \rho(\phi_a^\dagger)$

Kinetic > potential energy :  $\frac{1}{2} \dot{\phi}_a^2 > V(\phi_a)$

Kinetically dominated field evolves as  $a^{-6}$

Radiation evolves as  $a^{-4}$ .

**Ensures sufficient R.D.**

\*  $\phi_a$  refers to *after* instant preheating.  
 † We assume all the energy is instantaneously converted into radiation and thermalised

# Model Constraints

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Sub-Planckian:  $\phi < m_P$

DE observations:  $\rho(\phi_0) = \rho_{DE}$

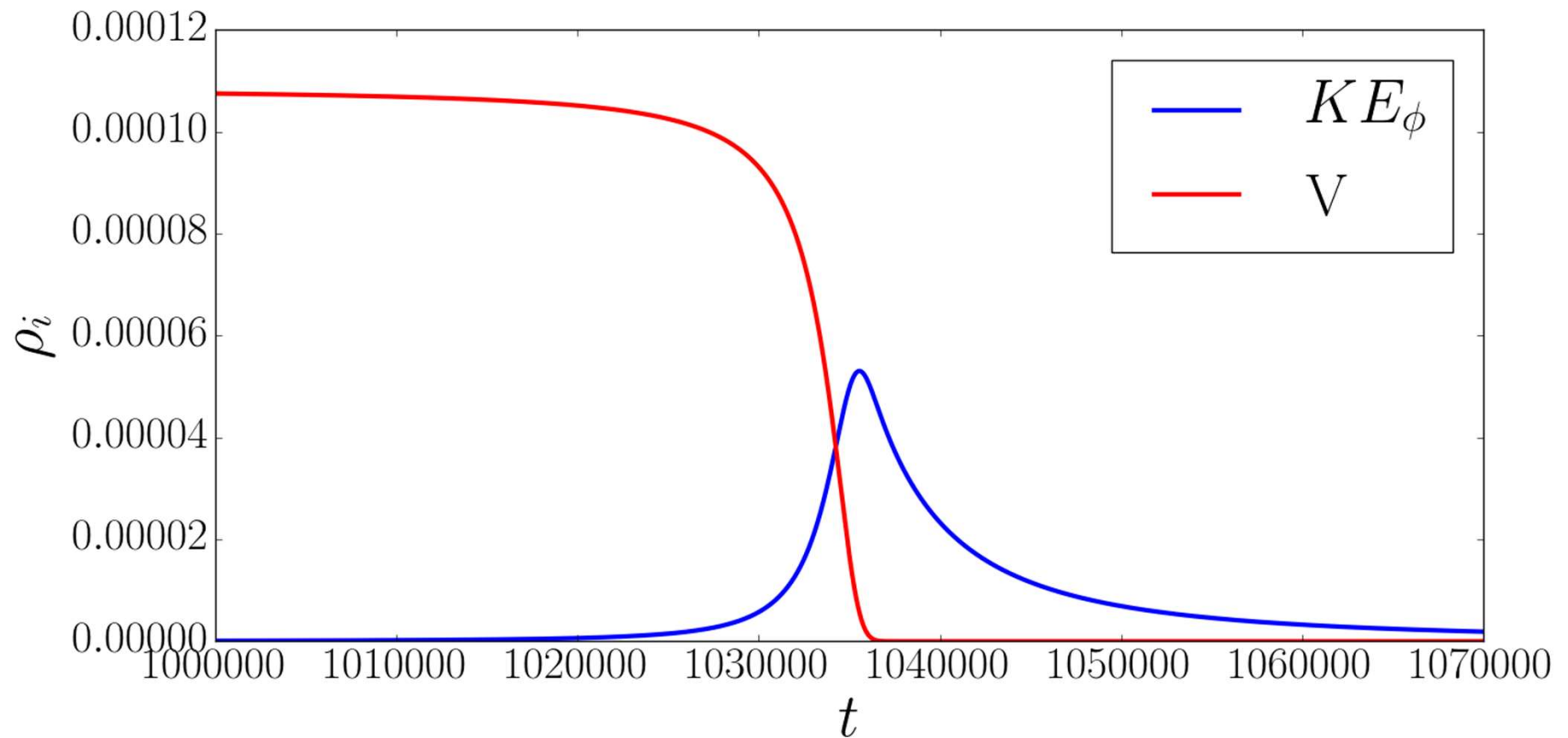
**Radiation Domination:**

$$\frac{1}{2} \rho(\phi_b^*) < \rho_\chi < \rho(\phi_b) - 2V(\phi_{IP})$$

\*  $\phi_b$  refers to  
*before* instant  
preheating.

$$\text{K. E.}(\phi) > 3V(\phi)$$


# Model Constraints



# Dark Energy Constraint

$$\rho(\phi_0) = \rho_{DE}$$


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- GB Coupling freezes time evolution of field:
- At very late times:

$$\frac{\phi_s}{m_P} = \frac{1}{q-2p} \ln \frac{2qV_0G_0}{3pm_P^2} \quad (1)$$

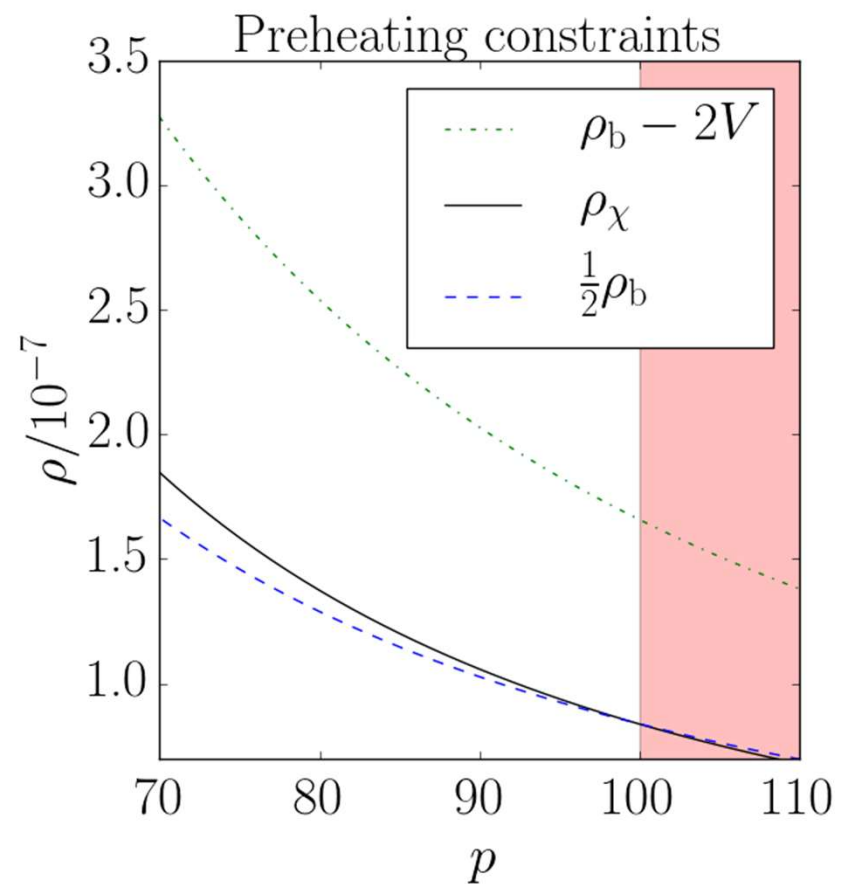
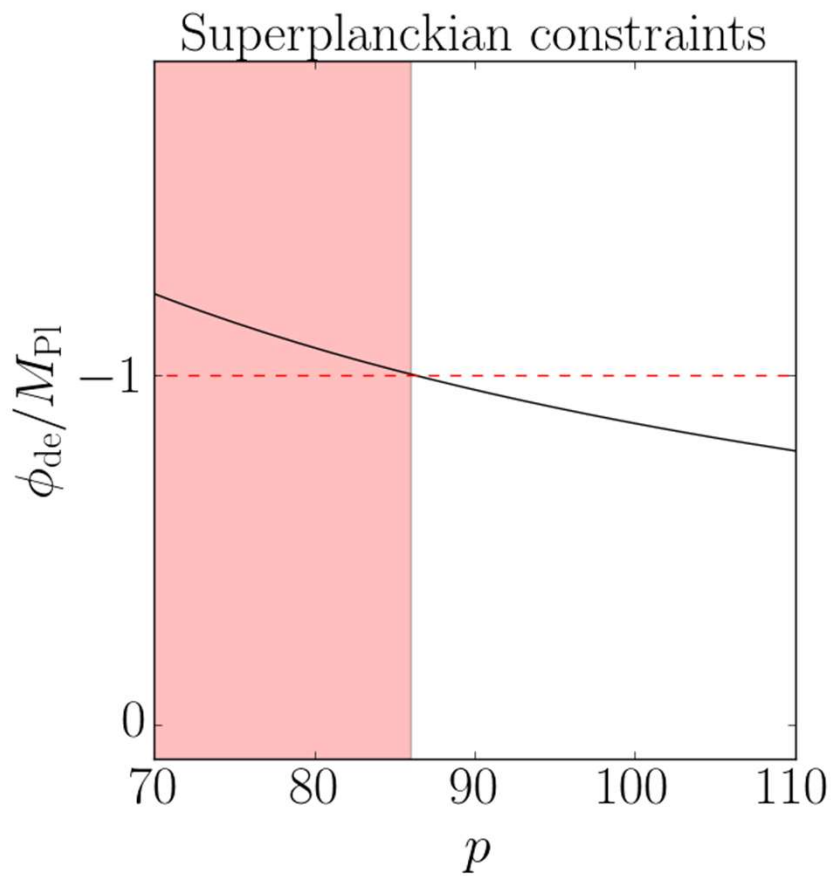
- In terms of the density parameters:

$$V + \frac{V_{,\phi}^2}{18H^2} + \left[ 3\Omega_\Lambda + \frac{2}{3}(7+3w)V_{,\phi}G_{,\phi} \right] m_P^2 H^2 + 2(1+3w)(13+3w)(m_P^2 G_{,\phi})^2 H^6 = 0 \quad (2)$$

- Inserting  $\Omega_\Lambda = 1$  into Eq.(2)  $\rightarrow$  we recover Eq.(1)
- Inserting  $\Omega_\Lambda = 0.7, \Omega_m = 0.3, H = H_0 \rightarrow$  we find  $\phi_0$
- We can use dark energy constraint to constrain parameters



# Results



$$q = 4p, G_0 m_P^2 = 1, g = 0.8$$

# Results


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- Natural choice of parameter values:

$$g < 1, G_0 > m_{\bar{P}}^{-2}, q > 2n$$

- Final constraints for *prototype* model:

$$50 \lesssim p \lesssim 350$$

- Large viable parameter space
  - Working model realisation applicable to any potential with two plateau regions
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# Theoretical Consideration

## - Fifth Force Problem

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- Mass dependent couplings to the standard model:

$$m_{eff}^2 = V'' \approx H$$

- Compton Wavelength  $\sim H^{-1}$  – we must include the interaction terms in the Lagrangian:

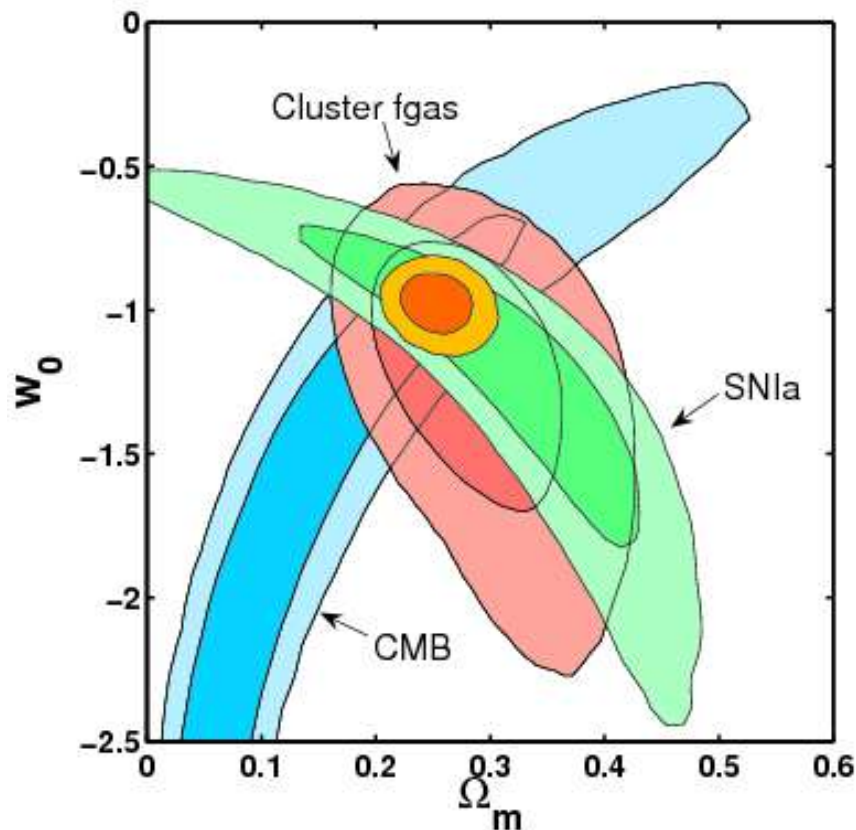
$$\sim \frac{\beta_i \phi}{m_p} \mathcal{L}_i$$

- To ensure suppression of these terms:

$$\beta_i \phi < m_p$$

- Motivation for our paper,  $\phi < m_p$ , avoids fine-tuning of  $\beta_i$ .
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# Distinguishing Between Cosmological Constant and Quintessence




## Experimental Observations of the Equation of State

- Cosmological constant,  $\omega = -1$
- Quintessential inflation, not always  $\omega = -1$ , time dependent.

# Gauss-Bonnet Coupled Quintessential Inflation - Summary

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- Coupling between Gauss-Bonnet term and scalar field maintains a **sub-Planckian field**.
  - Ensures the stability of the second plateau and **late-time acceleration**.
  - Inflationary observables in **agreement with Planck** results.
  - Parameter space **matching dark energy observations** today.
  - Quintessential Inflation is a **natural** (minimal and economical) alternative to  $\Lambda$ CDM
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# Thank you for listening, any questions?

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